



Impurity Densities from CXRS and Beam Emission Spectroscopy

R. Dux¹, B. Geiger¹, R. McDermott¹ and ASDEX Upgrade team¹

Experimental Validation of Ar CX Cross Sections

F. Guzman^{2,3}, M. Sertoli¹, R. McDermott¹, R. Guirlet³, L. Menchero^{1,3}, G. Modet⁴ and ASDEX Upgrade team¹

MPI für Plasmaphysik, EURATOM Assoziation, Garching bei München
ADAS-EU, University of Strathclyde
IRFM-CEA, Cadarache
Laboratoire Aimé Cotton, Université Paris XI

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Impurity densities from CXRS

- CX cross sections σ
- effective CX emission rate coefficients q
- D-density distribution in n=1 and n=2
 - depends on beam geometry, attenuation, excitation, distribution of species with full, half and third energy

$$\langle n_{Z+1} \rangle = \frac{4\pi}{h\nu} \frac{L_{CX}}{\sum_{j,k} \int n_{D,j,k} dl}$$

j=energy component, k=main quantum number

- Hα beam emission yields line-integrated densities in n=3
- thermal beam halo neutrals are produced by CX from beam neutrals onto thermal deuterons and produce a considerable fraction of the active impurity emission

$$\overline{\mathbf{A}^{+Z+1}} + \mathbf{D}(n = 1, 2..) \Longrightarrow \mathbf{A}^{+Z}(n, l) + \mathbf{D}^{+Z}(n, l) + \mathbf{D}^{+Z}(n, l) + \mathbf{D}^{+Z}(n, l)$$







Experimental Validation of Ar CX cross sections





- Ar used for radiation control and impurity transport experiments in fusion plasmas
- Ar charge exchange cross-sections from ORNL, ADAS scaling formula, and University Madrid differ by order of magnitude!
- Obtain experimental verification of correct calculation through comparison to X-ray data



Cross-sections differ by order of magnitude

- UAM (Madrid) data and ADAS universal formula cross-section order of magnitude larger than ORNL
- Differences are energy dependent
- AUG NBI can sample significant portion of energy spectrum





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Summary of Shots



Discharge	Line CER (NBI 1)	Line CHR (NBI 2)	Line CFR (NBI 1)	NBI 1/NBI 2 (KeV)
28263	522.4nm (ArXVIII)	522.4nm (ArXVIII)	661.2nm (ArXVI)	90/60KeV
28265	585.7nm (ArXVII))	585.7nm (ArXVII))	479.4nm (ArXVII)	90/60KeV
28268	541.2nm (ArXVI)	541.2nm (ArXVI)	522.4nm (ArXVIII)	75/50KeV
28269	427.6nm (ArXVIII)	427.6nm (ArXVIII)	541.2nm (ArXVI)	75/50KeV
28270	479.4nm (ArXVII)	479.4nm (ArXVII)	630.3nm (ArXVIII)	75/50KeV
28271	541.2nm (ArXVI)	541.2nm (ArXVI)	522.4nm (ArXVIII)	90/60KeV

 Almost complete energy scans for best Ar line per charge state

lon	Line	90KeV	75KeV	60KeV	50KeV
Ar XVIII	630.3nm n=17-16				Yes
Ar XVIII	522.4nm n=16-15	Yes		Yes	Yes
Ar XVIII	427.6nm n=15-14		No Sig		No Sig
Ar XVII	585.7nm n=16-15	Un Rel		Un Rel	
Ar XVII	479.4nm n=15-14		Yes	Yes	Yes
Ar XVI	661.2nm n=16-15			Yes	
Ar XVI	541.2nm n=15-14	Yes	Yes	Yes	Yes







- SXR density measurement:
 - Local experimental SXR emissivity (background subtracted & Abelinverted)
 - Electron density
 - Fractional abundance f_z
 - SXR filtered photon emissivity rate coefficients k_z

→ The denominator is evaluated using experimental n_e and T_e and calculating the ionization equilibrium assuming a standard set of transport coefficients





Ar Density from Soft X-Ray (SXR)

- SXR normalized emissivity (dashed = using local ionization eq.)
- Total Ar density (point with error bar = passive spectroscopic measurement of Ar¹⁶⁺ resonance lines)

 Ar density of 16+ → 18+ (fractional abundance includes transport)











something is wrong with the ionisation balance $\dots \rightarrow$ work in progress





Impurity Densities from CXRS and Beam Emission Spectroscopy



- 60keV beams from NBI I
- 25 lines-of-sight (LOS) for CXRS and 14 LOS for BES
- LOS aligned to centre of beam 3
- vertical separation ≈1.6cm
- Boron: n=7-6 transition at 494.7 nm
- H-Mode discharge with P_{NBI}=5 MW, P_{ECRH}=0.7 MW, medium density 6x10¹⁹m⁻³
- beam blips (200ms) active beam off for 100ms and replaced by beam of NBI II to have constant power

Cross-section of beam density and lines-of-sight (+ = BES, * = CXRS)





Beam Emission Spectrum

Doppler effect

$$\Delta \lambda_{dop} = \lambda_0 \frac{v}{c} \cos \vartheta$$

 motional Stark splitting due to electric field F in the rest frame of the beam atoms

$$\Delta \lambda_{\text{MSE}} = 2.76 \times 10^{-2} \, k \left| \vec{v} \times \vec{B} \right| \left[\frac{\text{MV}}{\text{m}} \right]$$
$$k = -4, -3, \dots, 4$$

 beam emission spectrum on top of halo emission (shifted Gaussian) – here for a case with high T_i









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IPP

setup CR models (cross-section data from Janev and ADAS) to calculate

- excited state population of D in beam
- beam stopping coefficients
- excited state population of D in halo
- halo production and stopping coefficients
- ratio of halo to beam particles





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D-densities along CXRS LOS from $D\alpha$

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D-densities along CXRS LOS from $\text{D}\alpha$

- halo radiance and density dominant
- good agreement with calculated beam and halo density distribution



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Boron Concentrations

radiance of impurity line

$$L_{cx} = c_{imp} \left[\sum_{j} L_{\alpha,j} \frac{q_{cx,j}}{q_{\alpha,j}} + L^{h}_{\alpha} \frac{\sum_{j} f_{j} q^{h}_{cx,j}}{\sum_{j} f_{j} q^{h}_{\alpha,j}} \right]$$

- boron signal mainly produced by beam species with full energy and halo
- contribution per halo neutral small but halo neutrals by far dominant species
- halo charge transfer only from excited atoms





Halo contribution for other popular lines





ratio of photons induced by halo N_{halo} to photons induced by beam N_{beam}

- Iargest for low energy beams due to large halo production by CX
- at low T only excited halo neutrals contribute at large T also charge transfer from ground state (mainly for He, less for heavier elements)





- Atomic data for calculation of excited state population of D are sufficiently good to get extra information on the beam attenuation from beam emission spectroscopy.
- A combined treatment of BES data and beam attenuation calculations will be the optimum solution for future analysis tools.
- CX excited impurity radiation from halo neutrals is an important contribution to the *active* CXRS signal.